# Physical and Magnetic Properties of Calcium-Substituted Li-Zn Ferrite

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(Submitted February 14, 2006)

The effect of Ca-substitution on the physical and magnetic properties of  $Li_{0.3-0.5x}Zn_{0.4}Ca_xFe_{2.3-0.5x}O_4$  ferrites (x = 0.0, 0.01, 0.02, 0.03, and 0.05), prepared by the standard ceramic method, has been studied. It is found that the saturation magnetization increases up to x = 0.01 and then it decreases. On the other hand, the initial permeability decreased while the Curie temperature remained almost constant with increasing x. The coercivity and remanence ratio increased with increasing x.

Keywords coercivity, electrical resistivity, initial permeability, magnetization, remanence, X-ray diffraction

## 1. Introduction

Li-Zn ferrites are low-cost materials and have important properties for technological applications. It is known that the intrinsic parameters of ferrites, such as magnetization, initial permeability, Curie temperature, and resistivity depend on the chemical composition, heat treatment, and type of additive or substituted ions (Ref 1). The rectangularity of B-H loop is also found to increase with increasing the defects that hinder the magnetic domain wall motion in the grains (Ref 2). The influence of additives, such as Na<sub>2</sub>O, CaO, Sb<sub>2</sub>O<sub>3</sub>, and ZrO<sub>2</sub> on the structure, magnetic and electrical properties of Li-Zn ferrites were previously reported (Ref 1, 3). This paper is devoted to study the effect of Ca<sup>2+</sup> substitution, not additives, on the physical and magnetic properties of Li-Zn ferrite.

## 2. Experimental Techniques

Ferrite samples with the chemical formula  $Li_{0.3-}$  $_{0.5x}Zn_{0.4}Ca_xFe_{2.3-0.5x}O_4$  (x = 0.0, 0.01, 0.02, 0.03, and 0.05) were prepared by the standard ceramic method. High purity oxides, 99.99%, of ZnO and Fe<sub>2</sub>O<sub>3</sub> with Li<sub>2</sub>CO<sub>3</sub> and CaCO<sub>3</sub> were mixed together according to their molecular weight. The mixture of each sample was ground to a very fine powder and then presintered at 800 °C for 6 h. The presintered mixture was ground again and pressed at room temperature, under a pressure

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of  $7 \times 10^8$  Pa, into tablet and toroidal forms. They were finally sintered at 1050 °C for 3 h and then slowly cooled to room temperature. X-ray diffraction patterns were performed using a diffractometer of type X'Pert Graphics and identified with Cu  $K_{\alpha}$  radiation. The theoretical x-ray density  $(d_x)$  of the samples was calculated using the formula  $(d_x = 8M/Na^3)$  where M is the molecular weight, N is Avogadro's number, and a is the lattice parameter. The density (d) of each composition was measured in bidistilled water using Archimedes' principle. The porosity percentage P (%) was calculated according to the relation  $P = 100(1 - d/d_x)$ %. The magnetization (M), at room temperature, was measured using the vibrating sample technique. The magnetizing field ranged from 0.0 up to 1.4 T. A toroidal sample of each composition was used as a transformer core for measuring the initial permeability  $\mu_i$  which was measured as a function of temperature at a constant frequency f = 10 KHz. The value of  $\mu_i$  was calculated using Poltinnikov's formula (Ref 4),  $V_s = K\mu_i$ , where  $V_s$  is the induced voltage in the secondary coil and  $K = \mu_0 n_p n_s I_p A \omega/L$ ;  $n_p$  and  $n_s$ are the number of turns of the primary and secondary coils, respectively,  $I_{\rm p}$  is the current in the primary coil, A is the cross-sectional area of the sample,  $\omega$  is the angular frequency, and L is the average path of the magnetic flux. Moreover, the coercive field and remanence were determined from the hysteresis loop. The tablet samples were used for the electrical resistivity measurements. The sample was inserted between two platinum electrodes where In-Hg was used as a contact material.

## 3. Results and Discussion

The x-ray diffraction patterns, (Fig. 1), showed that all investigated samples have single-cubic spinel phase. The d-spacing for each peak was recorded automatically and then the lattice parameter (a) was calculated from the relation

$$a = d_{hkl}(h^2 + k^2 + l^2)^{1/2}$$

The values of the lattice parameters obtained for each reflected plane are plotted against the function  $F(\theta)$  where  $F(\theta) = (\frac{1}{2})[(\cos^2\theta/\sin\theta) + (\cos^2\theta/\theta)]; \theta$  is the Bragg's angle.

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Straight lines were obtained and the accurate values of (*a*) were determined from the extrapolation of these lines to  $F(\theta) = 0$  (Ref 5). The variation of the lattice parameter (*a*) as a function of Ca-concentration (*x*) is shown in Fig. 2. It can be seen that the lattice parameter slightly increases linearly with the Ca-concentration, which obeys Vegard's law. Such an increase can be understood from the fact that we substitute Fe<sup>3+</sup> and Li<sup>1+</sup> ions by Ca<sup>2+</sup> ion which has a larger ionic radius ( $r_{Fe^{3+}} = 0.64$  Å,  $r_{Li^{1+}} = 0.68$  Å, and  $r_{Ca^{2+}} = 0.99$  Å). It is also noticed that the value of the lattice parameter (*a*) for the unsubstituted sample Li<sub>0.3</sub>Zn<sub>0.4</sub>Fe<sub>2.3</sub>O<sub>4</sub> (8.3817 Å) compares well with those reported earlier (8.378 Å (Ref 1, 6), 8.380 A (Ref 3)).

## 4. Magnetic Studies

#### 4.1 Magnetization

Figure 3 indicates the variation of magnetization M (emu/g) vs. the applied magnetic field H(T) at the room temperature. As a normal behavior, the magnetization increases with increasing the applied magnetic field and attains its saturation value for



Fig. 1 X-ray diffraction patterns of Li<sub>0.3-0.5x</sub>Zn<sub>0.4</sub>Ca<sub>x</sub>Fe<sub>2.3-0.5x</sub>O<sub>4</sub>



**Fig. 2** The variation of the lattice parameter a (Å) with Ca-concentration (x)

higher fields. The saturation magnetization  $M_s$  for each sample was determined by extrapolation of the magnetization curve to H = 0. The dependence of  $(M_s)$  on the Ca-concentration (x) is shown in Fig. 4. It is noticed that  $M_s$  increases rapidly with x and attains a maximum value at x = 0.01. Further increase of Ca-concentration leads  $M_s$  to decrease. Such a result could be discussed assuming the following cation distribution:

$$(Zn_{0.4}Ca_xFe_{0.6-x})[Li_{0.3-0.5x}Fe_{1.7+0.5x}]O_4$$

where the brackets () and [] denote A- and B-sites, respectively. Such cation distribution is based on the following facts:

- Li<sup>1+</sup> ions have a strong preference to occupy the B-site (Ref 7) while Zn<sup>2+</sup> ions have a strong preference to occupy the A-site (Ref 8).
- 2. As for Ca-ion distribution, it was reported that Ca<sup>2+</sup> ions strongly prefer to occupy the A-site for low Ca-concentration (Ref 6, 9). However, for high Ca-concentration, it was suggested that the Ca<sup>2+</sup> ions are either distributed between A- and B-sites (Ref 10) or reside at the grain boundaries (Ref 11). Thus, according to the assumed cation distribution (Ca ions occupy the A-sites) increasing



**Fig. 3** The magnetization (emu/g) vs. the applied magnetic field H (*T*)



Fig. 4 The variation of saturation magnetization  $M_{\rm s}$  (emu/g) with Ca-concentration (x)

of the Ca-concentration from x = 0 to x = 0.01 leads Fe<sup>3+</sup> content in B-site to increase and that in A-site to decrease. Hence the total magnetization  $(M_s = M_B - M_A)$  increases. For x > 0.01, more Fe<sup>3+</sup> ions migrate to B-site causing the B-B interaction to increase and hence the canting angle  $(\theta_{YK})$  establishes. It is expected that  $(\theta_{YK})$  increases with increasing Ca-concentration. Such an increase in  $(\theta_{YK})$  leads to a decrease of  $M_s$  according to the equation (Ref 12)  $M_s = M_B \cos \theta_{YK} - M_A$ , where  $M_A$  and  $M_B$  are the magnetizations of A- and B-sites, respectively. Thus our results support the occupation of Ca ions for A-sites at low Ca-concentration.

#### 4.2 Initial Permeability

Figure 5 shows the variation of the initial permeability  $\mu_i$  with temperature for the series  $\text{Li}_{0.3-0.5x}\text{Zn}_{0.4}\text{Ca}_x\text{Fe}_{2.3-0.5x}\text{O}_4$  (x = 0.0, 0.01, 0.02, 0.03, and 0.05). It is seen that the initial permeability increases with temperature up to Curie temperature  $T_c$ . Near  $T_c$  there is a sharp drop in  $\mu_i$ . This result could be explained according to the Globus relation (Ref 13), which is given by

 $\mu_i \propto (M_s^2 D/K_1^{1/2})$ 

where *D* is the average grain size and  $K_1$  is the anisotropy constant. According to the Globus relation, the dependence of  $\mu_i$  on temperature results from the temperature dependence of  $M_s$  and  $K_1$ . It is known that the anisotropy field usually decreases much faster with temperature than  $M_s$  (Ref 14) which leads to an increase in  $\mu_i$ . Figure 5 shows also that the rate of increase of  $\mu_i$  with temperature decreases by increasing Cacontent. This may lead us to conclude that the Ca<sup>2+</sup> ions improve the thermal stability of  $\mu_i$  for the investigated samples (Ref 14) by compensating changes in  $M_s$  and  $K_1$  with temperature. The compensating mechanism is still under investigation.

Furthermore, from Fig. 5 one can conclude that  $\mu_i$ , at room temperature, decreases with increasing Ca-concentration. This is illustrated in Fig. 6 in addition to the variation of the porosity with Ca content. It is obvious that as the porosity increases,  $\mu_i$ 



Fig. 5 The temperature dependence of the initial permeability  $\mu_i$  for different Ca-concentrations (*x*)

decreases with increasing Ca-concentration. The decrease of  $\mu_i$  with increasing *x* could be discussed using Globus relation as follows.

According to the proposed cation distribution, the amount of iron in B-site increases, which in turn increases the amount of  $Fe^{2+}$  in B-site. According to the one ion model, the anisotropy constant  $K_1$  has a direct relation with the amount of  $Fe^{2+}$ . Therefore, as Ca-concentration increases,  $K_1$  also increases. Furthermore, previous studies (Ref 1) have shown that doping Li-Zn ferrite with CaO decreases the grain size *D*. These two factors and the increase in porosity lead to the decrease of  $\mu_i$  with increasing *x*. Moreover, the decrease of  $M_s$  in the range  $0.01 < \times \le 0.05$ , may contribute in decreasing  $\mu_i$ .

# 5. Curie Temperature

The values of Curie temperature ( $T_c$ ) were determined from the extrapolation of the linear part at the sudden decrease of  $\mu_i$ of Fig. 5 for all investigated samples. The values of  $T_c$  and the rate of decrease of  $\mu_i$  (slope of the linear part) are given in Table 1.

Table 1 illustrates that  $T_c$  is almost constant (628±5 K). The constancy of  $T_c$  is discussed as follows. It is known that, in ferrites,  $T_c$  depends mainly on the A-B interaction, which in turn is affected by the following factors:

 The magnitude and the concentration of magnetic moments existed in A- and B-sites. In our investigated system, the only moments existed those of Fe<sup>3+</sup> ion. Moreover, the relative concentration of the moments on



**Fig. 6** The variation of the initial permeability  $\mu_i$ , at room temperature, and porosity *P* (%) with Ca-concentration (*x*)

Table 1	The values	of $T_{\rm c}$ (K	) and	rate of	decrease	of $\mu_i$
with Ca-o	concentratio	n				

Ca-Conc. (x)	<i>T</i> <sub>c</sub> (K)	Rate of decrease of $\mu_i$
0.0	628	-15.85
0.01	628	-9.25
0.02	623	-5.56
0.03	630	-3.97
0.05	633	-2.07

A- and B-sites  $(M_A/M_B)$  is almost constant, where for unsubstituted sample  $(M_A/M_B) = 0.353$ , and for the highest substituted sample  $(M_A/M_B) = 0.319$ .

2. The distance between the A- and the B-sites where Fig. 2 shows that the change of lattice parameter with Ca-concentration, 0.025 Å/Ca content, is very small to affect  $T_{\rm c.0}$ 

Thus, according to the above factors,  $T_c$  tends to be almost constant which we have found experimentally.

Table 1 indicates also that the slope of the linear part of  $\mu_i(T)$  curve, at the sudden decrease of  $\mu_i$ , decreases with increasing Ca<sup>2+</sup> concentration. It was reported that the value of this slope gives a good indication about the sample homogeneity (Ref 15), the higher the slope the higher the homogeneity. Thus one concludes that the homogeneity decreases with increasing Ca-concentration. Such a result could be attributed to the increase of the porosity with Ca-concentration.

## 6. Coercivity and Squareness

The variations of the coercive field  $H_c$  and the remanence ratio  $(B_r/B_s)$ , which expresses the squareness of the hysteresis loop, for all Ca-concentrations (*x*) at a field H = 340 A/m are given in Table 2.

One can observe that  $H_c$  increases with increasing x. Similar results were reported by Rezlescu et al. (Ref 16) for Ni-Zn ferrite doped with CaO. Table 2 shows also that the ratio  $(B_r/B_s)$ increases with increasing Ca-concentration (x). These results are in agreement with that of the porosity. As the porosity increases, then higher field is needed to push the domain wall, i.e.,  $H_c$  increases. Moreover, the behavior of  $M_s$  may play an indirect role in increasing  $H_c$ , specially for x > 0.01, through Brown's relation (Ref 17) which is given by

 $H_{\rm c} \geq (2K_1/\mu_0 M_{\rm s})$ 

Furthermore, it is known that the remanence  $B_r$  is directly proportional to the porosity (Ref 18). Moreover, for  $x \le 0.01$ , the porosity increases by a factor of 67% while  $M_s$  increases only by a factor of 11.5%. This leads to an increase of  $B_r/B_s$ . In addition, for the range  $0.01 < x \le 0.05$ , the porosity continues to increase while  $M_s$  decreases. This causes more increase of  $B_r/B_s$ .

#### 6.1 Resistivity at Room Temperature

The dc resistivity ( $\rho$ ) of all samples, at room temperature, are given in Table 3.

The table indicates that in the range  $0 < \times \le 0.02$ , the dc resistivity ( $\rho$  decreases with increasing Ca-concentration (*x*), while  $\rho$  increases again for x > 0.02. This result could be explained as follows:

It was assumed above that the  $Ca^{2+}$  ions occupy the A-site for Ca-concentration up to x = 0.01. Hence, some  $Fe^{3+}$  ions migrate from A- to the B-site. The migration of  $Fe^{3+}$  is accompanied by the formation of  $Fe^{2+}$  ions in B-site (Ref 19). Thus the ratio of  $Fe^{2+}/Fe^{3+}$  ions in B-site increases. Thus the probability of hopping between  $Fe^{2+}$  and  $Fe^{3+}$  increases and hence, the resistivity decreases. For higher Ca-concentration, the effect of the porosity becomes the dominant factor and the resistivity increases.

Table 2 The values of  $H_c$  (A/m) and  $(B_r/B_s)$  with Caconcentration (x)

Ca-Conc. (x)	$H_{\rm c}$ (A/m)	$(B_{\rm r}/B_{\rm s})$	
0.0	92	0.17	
0.01	100	0.22	
0.02	111	0.32	
0.03	120	0.33	
0.05	143	0.44	

Table 3 The variations of  $\rho$  ( $\Omega$  cm) with Ca-concentration

Ca-Conc. (x)	$\rho (\Omega \text{ cm}) \times 10^7$	
0.0	3.6	
0.01	0.1	
0.02	0.08	
0.03	2.1	
0.05	30.4	

## 7. Conclusion

Single spinel phases of Ca-substituted Li-Zn ferrite are studied. The Ca-substitution slightly increased the lattice parameter, the coercivity and the squareness of the samples. The magnetization increased for small Ca-concentration, x = 0.01, and then decreased for higher Ca content. The homogeneity of the samples and initial permeability decreased with increasing Ca-concentration while Curie temperatures are found to be almost constant.

#### References

- N. Rezlescu and E. Rezlescu, The Influence of Additives on the Material Parameters of Li-Zn Ferrites, *Phys. Stat. Sol. (a)*, 1995, 147, p 553–562 (in English)
- H.T. Kim and H.B. Im, Effects of Bi<sub>2</sub>O<sub>3</sub> and Nb<sub>2</sub>O<sub>5</sub> on the Magnetic Properties of Ni-Zn Ferrites and Lithium Ferrites, *IEEE Trans. Magnet.*, 1982, **18**(6), p 1541–1543 (in English)
- E. Bermejo, T. le Mercier, and M. Quarton, Microstructure and Physicochemical Studies of Pure and Zinc-Substituted Lithium Ferrites Sintered Above 1000°C, *J. Am. Ceram. Soc.*, 1995, 78(2), p 365–368 (in English)
- S.A.J. Poltinnikov, Some Magnetic Properties of Nickel-Cadmium Ferrite, Sov. Phys. Solid Stat., 1966, 8, p 1144–1149 (in English)
- B.D. Cullity, *Elements of X-ray Diffraction*, Second printing, Addison-Wesley Publishing Company, Inc, USA, 1959
- N. Rezlescu and E. Rezlescu, Effects of Addition of Na<sub>2</sub>O, Sb<sub>2</sub>O<sub>3</sub>, CaO, and ZrO<sub>2</sub> on the Properties of Lithium Zinc Ferrite, *J. Am. Ceram. Soc.*, 1996, **79**(8), p 2105–2108 (in English)
- G.O. White and C.E. Patton, Magnetic Properties of Lithium Ferrite Microwave Materials, *J. Magn. Magn. Mater.*, 1978, 9, p 299–317 (in English)
- Z.C. Xu, Magnetic Anisotropy and Mössbauer Spectra in Disordered Lithium-Zinc Ferrites, J. Appl. Phys., 2003, 93(8), p 4746–4749 (in English)
- R.V. Upadhyay, G.J. Baldha, and R.G. Kulkarni, Mössbauer study of the Ca<sub>x</sub>Co<sub>1-x</sub>Fe<sub>2</sub>O<sub>4</sub> system, *J. Magn. Magn. Mater.*, 1986, **61**, p 109– 113 (in English)
- S.D. Chhaya, M.P. Pandya, M.C. Chhantbar, K.B. Modi, G.J. Baldha, and H.H. Joshi, Study of Substitution Limit, Structural, Bulk Magnetic

and Electrical Properties of Ca<sup>2+</sup> Substituted Magnesium Ferrite, J. Alloys Comp., 2004, **377**, p 155–161 (in English)

- A.Y. Lipare, P.N. Vasambekar, and A.S. Vaingankar, Ac Susceptibility Study of CaCl<sub>2</sub> Doped Copper–Zinc Ferrite System, *Bull. Mater. Sci.*, 2003, 26(5), p 493–497 (in English)
- S.S. Suryawanshi, V.V. Deshpande, U.B. Deshmukh, S.M. Kabur, N.D. Chaudhari, and S.R. Sawant, XRD Analysis and Bulk Magnetic Properties of Al<sup>3+</sup> Substituted Cu-Cd Ferrites, *Mater. Chem. Phys.*, 1999, **59**, p 199–203 (in English)
- G.C. Jain, B.K. Das, R.S. Khanduja, and S.C. Gupta, Effect of Intragranular Porosity of Initial Permeability and Coercive Force in a Manganese Zinc Ferrite, *J. Mater. Sci.*, 1976, **11**, p 1335–1338 (in English)
- A.M. Sankpal, S.V. Kakatkar, N.D. Chaudhari, R.S. Patil, S.R. Sawant, and S.S. Suryavanshi, Initial Permeability Studies on Al<sup>3+</sup> and Cr<sup>3+</sup>

Substituted Ni-Zn Ferrites, J. Mater. Sci. Mater. Electron., 1998, 9, p 173-179 (in English)

- A. Globus, H. Pascard, and V.J. Cagan, Distance Between Magnetic Ions and Fundamental Properties in Ferrites, *J. De Phys. Colloque*, 1977, 38(4), p C1-163–C1168 (in English)
- E. Rezlescu, N. Rezlescu, C. Pasnicu, M.L. Craus, and D.D. Popa, The Influence of Additives on the Properties of Ni-Zn Ferrite Used in Magnetic Heads, J. Magn. Magn. Mater., 1992, 117, p 448–454 (in English)
- J.M.D. Coey, Rare Earth Permanent Magnetism. 1st ed., John Wiley and Sons, New York, 1996, 220
- R.M. Bozorth, *Ferromagnetism*. 7th ed., Van Nostrand Company, Inc, Princeton, New Jersey, 1963, 824
- N. Miyata, Electric Conduction of Ferrites Containing Fe<sup>2+</sup>-ions, J. Phys. Soc. (Jpn), 1961, 16, p 206–208 (in English)